CENTRE FOR MICR O- AND NANOTECHNOLOGY



Kylmälaboratorio Lågtemperaturlaboratoriet Low Temperature Laboratory

Low Temperature Laboratory in Micronova

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Jukka Pekola

Low Temperature Laboratory in Micronova

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- 1. Electronic micro-refrigeration and cold electron Josephson transistor
- 2. Single Cooper pair pumping

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island

Probine

unction









Charge and Flux Controlled Pumping of Cooper pairs

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- 1. Motivation
- 2. Principle of the device
- 3. Experiments





Why to pump charges?

Towards current standard

"Quantum triangle"



Normal single-electron pump



M. W. Keller, J. M. Martinis, N. M. Zimmerman, and A. H. Steinbach, Appl. Phys. Lett. **69**, 1804 (1996). High accuracy but slow: I < 10 pA





Why to pump Cooper pairs?

Cooper pair pumps could possibly be faster than normal electron pumps



Charge transport adiabatic up to

$$f_{\rm LZ} \sim \frac{E_{\rm J}^2}{\hbar E_{\rm C}}$$

Interesting also in terms of quantum coherent effects in small Josephson junction circuits CENTRE FOR MICRO- AND NANOTECHNOLOGY



But: there are errors due to large $E_{\rm J}$

Perfectly phase-biased adiabatic CPP





$$Q_P = 2\hbar \Im m \left[\sum_{n \neq m} \oint \frac{(\hat{I}_l)_{mn}}{E_m - E_n} \langle n | \partial_{\vec{q}} m \rangle \cdot d\vec{q} \right]$$

 $Q_{\rm P}/(2e) \simeq 1 - 9E_{\rm J}/E_{\rm C}\cos\varphi$

This adds to "normal" supercurrent $I_{\rm S} \propto E_{\rm J} \sin \varphi$

J. P. et al., PRB 60, 9931 (1999); R. Fazio et al., PRB 68, 054510 (2003).





The first three-junction CPP



L. Geerligs et al., Z. Phys. B: Condensed Matter 85, 349 (1991).

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Can be generalized to pump 2Ne per cycle. (N = 1, 2, ...?)

ω



Predicted accuracy of the (ideal) device



To study also the non-adiabaticity errors, we obtained these results by solving Schrödinger eq. and integrating in time, not by adiabatic approximation.



Influence of residual Josephson coupling and offset charge



Phase noise may average part of the $E_{\rm J}^{\rm res}$ -error out.



Potential sources of error

- 1. non-adiabaticity
- 2. non-ideal suppression of $E_{\rm J}$
- 3. environmental impedance
- 4. background charge noise
- 5. quasiparticles





The measured device



Input coils



Junctions

Gate line

SQUID loops







General *IV* curves, pumping

3 MHz, 4...34 pairs / cycle pumped







Quantitative comparison to I = N2ef



30 pairs per cycle at 4 MHz with accuracy of 1 % achieved



Improvements in the future

- 1. Better temporal suppression of $E_{\rm J}$ by threejunction SQUID or SQUID array
- 2. Better designed environment to suppress supercurrent leakage (fixed voltage bias)
- 3. Get rid of quasiparticles (by gap engineering?)
- 4. Higher speed by
- (a) increasing $E_{\rm J}$ (by lowering junction resistance or ultimately by using Nb junctions)
- (b) pulse optimisation to avoid non-adiabaticity

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Hysteretic DC-SQUIDs with low critical current as read-out devices in quantum circuits



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DC SQUID or Josephson junction as an ammeter: applications

Qubits and coherent Cooper pair pumps – **avoid dissipation**



"Classical" measurements (e.g. noise and FCS) can tolerate some dissipation



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Increasing sensitivity of Josephson junction ammeters/threshold detectors

$$\Delta I \propto I_{\rm C}^{3/5}/C^{2/5}$$

Best resolution when both $E_{\rm J}$ ($I_{\rm C}$) and $E_{\rm C}$ (1/*C*) are small. But how far can we go?











Phase diagram of small Josephson junctions and SQUIDs







THE END

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ENERGY LEVEL MODEL and DISSIPATION

